

Determination of α_s using jet rates at LEP with the OPAL detector

The OPAL Collaboration

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Abstract. Hadronic events produced in e^+e^- collisions by the LEP collider and recorded by the OPAL detector were used to form distributions based on the number of reconstructed jets. The data were collected between 1995 and 2000 and correspond to energies of 91 GeV, 130–136 GeV and 161–209 GeV. The jet rates were determined using four different jet-finding algorithms (Cone, JADE, Durham and Cambridge). The differential two-jet rate and the average jet rate with the Durham and Cambridge algorithms were used to measure α_s in the LEP energy range by fitting an expression in which $\mathcal{O}(\alpha_s^2)$ calculations were matched to a NLLA prediction and fitted to the data. Combining the measurements at different centre-of-mass energies, the value of $\alpha_s(M_Z)$ was determined to be

$$\alpha_s(M_Z) = 0.1177 \pm 0.0006 (\text{stat.}) \pm 0.0012 (\text{expt.}) \pm 0.0010 (\text{had.}) \pm 0.0032 (\text{theo.}).$$

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1 Introduction

In the Standard Model of elementary particle interactions, the strong interaction is described by the theory of Quantum Chromodynamics (QCD), and depends on just one fundamental parameter, the strong coupling α_s . The value of α_s is expected to depend on the energy scale of the interaction. It is therefore an important test of the theory to determine the value of α_s experimentally at as many different energies as possible. It is also important to use as many different techniques as possible, as different measurements of α_s are sensitive to different theoretical and hadronization variations.

Indeed, many methods have already been employed to evaluate α_s [1]. At very low energies the value of α_s can be measured using the hadronic decays of the τ lepton and heavy quarkonia. Low energy determinations are also available using scaling violations and sum rules from deep inelastic scattering experiments. Higher energy determinations of α_s come from collider experiments (e^+e^- , pp , $p\bar{p}$ or ep) using properties of the created hadron system which are explicitly dependent on the value of $\alpha_s(Q)$, where Q corresponds to the energy scale at which the interaction takes place¹.

During the LEP1.5 ($\sqrt{s} \sim 133$ GeV) and LEP2 (above W^+W^- threshold) operational phases of the Large

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¹ For e^+e^- collisions $Q = \sqrt{s}$.

Electron-Positron collider at CERN, events were recorded with centre-of-mass collision energies ranging from 91 GeV to 209 GeV. Events of the form $e^+e^- \rightarrow \text{hadrons}$ can be used to determine distributions based on the ensemble of final state hadrons (event shapes) or on the ensemble of jets (jet rates). Previous results by OPAL for an α_s determination based on event shapes and jet rates using the Z dataset collected during the LEP1 phase can be found in [2]. Determinations of α_s from LEP1.5 and LEP2 datasets up to 189 GeV have already been reported by OPAL based on event shape distributions [3–5] and on jet rates [6]. Another OPAL paper [7] uses the same data that have been presented here to measure event shapes.

For the analysis presented in this paper we used data collected during the LEP1.5 and LEP2 phases to construct jet rate distributions using several jet clustering algorithms. The differential two-jet rate, D_2 , and the average jet rate, $\langle N \rangle$, were used to determine values of $\alpha_s(\sqrt{s})$ at the four combined centre-of-mass energies composed of data within the LEP1.5 and LEP2 datasets. Theoretical predictions were fitted to these distributions to extract the value of $\alpha_s(\sqrt{s})$.

The paper is organized as follows. Section 2 contains a brief description of the OPAL detector. A summary of the data and the Monte Carlo samples used in the analysis is given in Sect. 3. In Sect. 4, we define the jet rate distributions. The methods used to select signal events and reject backgrounds are presented in Sect. 5. The variations used for systematic studies are detailed in Sect. 6. Finally, the results of this analysis are given in Sect. 7, followed by a conclusion and summary in Sect. 8.

2 The OPAL experiment

A full description of the OPAL detector can be found in [8]. The critical components of the detector in the identification of jets were the central tracking chambers, which were used to reconstruct charged particles, and the electromagnetic calorimeters, which measured the total energy deposited by electrons and photons.

The tracking chambers were located inside a solenoidal magnet which provided a 0.435 T axial magnetic field along the beam axis. The main component of the tracking system was a large-volume jet chamber, which was approximately 4.0 m long with an outer radius of 1.85 m. The jet chamber was separated into 24 sectors, each with a radial plane of 159 sense wires separated by 1 cm. The momenta of tracks in the $x-y$ plane² were measured with a precision parametrized by $\sigma_p/p = \sqrt{0.02^2 + (0.0015 \cdot p[\text{GeV}/c])^2}$.

The calorimetry systems were outside the solenoidal magnet. The electromagnetic calorimeter was composed of 11704 lead glass blocks in the barrel and endcap regions, representing about 25 radiation lengths in the barrel and

more than 22 in the endcap. The iron sampling hadron calorimeter was located just outside the electromagnetic calorimeter, and provided the stopping power to contain most hadronic showers. Luminosity was determined using small-angle Bhabha events detected in the forward detectors and silicon-tungsten calorimeter [9].

After an event was triggered [10], data were collected from the subdetectors and processed by the OPAL data acquisition system [11]. The raw event data were transferred to a farm of computer processors where the events were fully reconstructed and written to tape for offline analysis.

3 Data and Monte Carlo samples

The data used in this analysis were collected by OPAL between 1995 and 2000 and correspond to integrated luminosities of 14.7 pb⁻¹ of data taken with centre-of-mass energy 91 GeV, 11.3 pb⁻¹ of LEP1.5 data with centre-of-mass energies between 130 GeV and 136 GeV and 707.4 pb⁻¹ of LEP2 data with centre-of-mass energies ranging from 161 to 209 GeV. The 91 GeV data, known as Z-calibration data, were primarily collected for calibrating parameters used in the OPAL reconstruction algorithms. This M_Z sample had the same detector configuration as the other centre-of-mass energy points. The exact breakdown of the centre-of-mass energies together with the respective luminosities and numbers of selected events are given in Table 1. The thirteen points in Table 1 represent the main samples of the spread of energies in the LEP1.5 and LEP2 data.

The data were combined into four datasets. The LEP1.5 data provided a single energy point at an event-weighted centre-of-mass energy of 133 GeV, while the LEP2 data were split into two energy points, one with an event-weighted centre-of-mass energy of 177 GeV using data in the range 161–185 GeV (with a total integrated luminosity of 78.1 pb⁻¹) and another at 197 GeV (with a total integrated luminosity of 628.3 pb⁻¹) using data in the range 188–209 GeV. Together with the Z-calibration data this provided for a determination of α_s at four centre-of-mass energies.

A number of Monte Carlo samples were created to correct for detector acceptance and resolution effects, to correct for hadronization effects and to estimate the contribution of background processes. These Monte Carlo samples were produced using a full simulation of the detector [12], followed by the same reconstruction and selection algorithms applied to the real data, and are referred to as “detector-level” samples. Other samples without the full detector simulation are discussed in Sect. 5.2.

PYTHIA 6.150 [13] was used to provide the default Monte Carlo samples (for the process $e^+e^- \rightarrow Z/\gamma^* \rightarrow q\bar{q} \rightarrow \text{hadrons}$) which were used to correct the high energy datasets. The Z-calibration dataset was corrected using JETSET 7.408 [14]. The use of JETSET for the lower energy data is a matter of convenience only, and not due to any inconsistencies in PYTHIA at this energy. Any differences between the two generators is expected to be negligible. Hadronization corrections were evaluated by comparing

² The right-handed OPAL coordinate system is defined so that z is the coordinate parallel to the e^- beam direction and the x axis points to the centre of the LEP ring, r is the distance normal to the z axis, θ is the polar angle with respect to the z axis and ϕ is the azimuthal angle with respect to the x axis.

Table 1. Integrated luminosity and the total number of preselected events for all samples. Included in this table are the Monte Carlo samples used to correct and make comparisons with the data. The 91 GeV dataset corresponds to the data collected during the Z-calibration runs

| Year | Energy (GeV) | | | Integrated Luminosity (pb^{-1}) | Number of Events | | |
|------------|--------------|-------------|-------|--|------------------|-------------------|--------|
| | Nominal | Range | Mean | | Data | JETSET/ PYTHIA | HERWIG |
| 1996–2000 | 91 | 91.0–91.5 | 91.3 | 14.7 | 426194 | 459k | 443k |
| 1995, 1997 | 130 | 130.0–130.3 | 130.1 | 5.3 | 1628 | 50k | 50k |
| 1995, 1997 | 136 | 135.7–136.2 | 136.1 | 6.0 | 1504 | 50k | 50k |
| 1996 | 161 | 161.2–161.6 | 161.3 | 10.1 | 1369 | 100k | 80k |
| 1996 | 172 | 170.2–172.5 | 172.1 | 10.4 | 1285 | 100k | 80k |
| 1997 | 183 | 180.8–184.2 | 182.7 | 57.7 | 6027 | 100k | 100k |
| 1998 | 189 | 188.3–189.1 | 188.6 | 185.2 | 18559 | 500k | 100k |
| 1999 | 192 | 191.4–192.1 | 191.6 | 29.5 | 2866 | 200k | 100k |
| 1999 | 196 | 195.4–196.1 | 195.5 | 76.7 | 7076 | 200k | 100k |
| 1999, 2000 | 200 | 199.1–200.2 | 199.5 | 79.3 | 6676 | 200k | 100k |
| 1999, 2000 | 202 | 201.3–202.1 | 201.6 | 37.8 | 3225 | 200k | 100k |
| 2000 | 205 | 202.5–205.5 | 204.9 | 82.0 | 6721 | 200k | 100k |
| 2000 | 207 | 205.5–208.9 | 206.6 | 138.8 | 10987 | 375k | 100k |

results with an alternative Monte Carlo sample, HERWIG 6.2 [15] which uses the cluster model of hadronization. This was compared with the string model of hadronization in PYTHIA. The parameters which were involved in the Monte Carlo simulation, both for JETSET/PYTHIA and HERWIG, were tuned to OPAL data collected at the Z peak, including global event shapes, particle multiplicities and fragmentation functions [16, 17]. The generation of the initial quark-antiquark pair for each Monte Carlo sample was implemented at LEP2 using the *KC2f* 4.13 event generator [18], which has an improved description of photon production in the initial and final states with respect to the one currently implemented in the PYTHIA generator. The available detector-level Monte Carlo samples are listed in Table 1.

Above the W^+W^- production threshold (161 GeV), the main background was expected to come from four-fermion events ($e^+e^- \rightarrow W^+W^- \rightarrow 4f$), in particular those events in which two or all four of the fermions were quarks. The contribution of these backgrounds in data was estimated using Monte Carlo samples generated using KORALW 1.42 [19] (for $q\bar{q}q'\bar{q}'$ and $q\bar{q}\ell\bar{\ell}'$ where $\ell = e, \mu, \tau, \nu$ but $\ell\bar{\ell} \neq e^+e^-$) and grc4f 2.1 [20] (for $eeq\bar{q}$). Grc4f 2.1 was used to generate all the expected four-fermion background samples for the 161 and 172 GeV data. The background distributions were normalized to the luminosity of the dataset and subsequently subtracted from the measured distributions. The LEP1.5 energies were well below the W^+W^- and ZZ production thresholds [21] and were therefore expected to have no significant four-fermion backgrounds. The total expected background contribution from “four-fermion” $e^+e^- \rightarrow q\bar{q}q\bar{q}$ events is 1.2% of the combined LEP1.5 data sample and it was neglected in the analysis.

4 Jet rate distributions

Jets were formed from the final state objects by applying jet clustering algorithms. These algorithms use the kinematic and spatial (geometric) properties of the individual objects in order to classify them as belonging to a specific jet. We used here the Durham [22], Cambridge [23], JADE [24] and the R and ε variants of the Cone [25] jet clustering algorithms.

The Durham and Cambridge algorithms construct a test variable built from the energy and angular separation between two particles,

$$y_{ij} = \frac{2 \min\{E_i^2, E_j^2\}(1 - \cos \theta_{ij})}{E_{\text{vis}}^2}$$

where E_i is the energy of particle i , θ_{ij} the angle between the particle i and j and E_{vis} is the total visible energy in the event. The pair that produces the smallest value of y_{ij} is chosen first. The value of this test variable is compared to a predefined parameter, y_{cut} , called the jet resolution parameter. If the test variable is smaller than y_{cut} particles i and j are merged into a pseudo-particle. Merging means that the momenta of particles i and j are removed from the set of momenta and the sum of their four-momenta is added to the set of momenta. After the merging, the clustering starts again using the momentum set and it continues until all test variables become larger than y_{cut} . After the clustering stops, all remaining (pseudo-) particles are classified as jets.

The Cambridge algorithm differs slightly from Durham in its implementation. In the Cambridge algorithm particles are paired, starting with the smallest y_{ij} , where the

variable $v_{ij} = 2(1 - \cos \theta_{ij})$. The standard test variable is then constructed and compared to the jet resolution parameter, y_{cut} . The procedure followed is then identical to that of the Durham algorithm, except that Cambridge freezes out soft jets by accepting only the lowest energy (pseudo-)particle as the jet when $y_{ij} > y_{\text{cut}}$. The number of jets reconstructed in the event, using Durham or Cambridge, is therefore a function of the jet resolution parameter. The JADE algorithm follows the same procedure as the Durham algorithm; however, it uses the scaled invariant mass, $y_{ij} = 2E_i E_j (1 - \cos \theta_{ij}) / E_{\text{vis}}^2$, of particles i and j as the test variable.

In the Cone jet finding algorithm, a jet is defined as a set of particles whose three-momentum vectors lie inside a cone of half angle R , where the direction of the sum of their three-momentum vectors defines the cone axis. In addition, the total energy of the particles assigned to a jet is required to exceed some minimum value ε . Typical values are $R = 0.7$ rad and $\varepsilon = 7$ GeV for jets in e^+e^- annihilation at LEP 1 energies. When analysing events at the detector level, we replaced ε by $\varepsilon' = \varepsilon \cdot E_{\text{vis}} / \sqrt{s}$ to compensate for the incomplete detection of the energy of the event. In our studies, the jet rate was computed at fixed $\varepsilon = 7$ GeV as R was varied, and at fixed $R = 0.7$ as ε was varied. The former is sensitive to the angular structure of jets, and the latter to their energy distribution.

The fraction of multihadronic events in a given sample that are classified as containing n jets for a given value of the jet resolution parameter (y_{cut} , R or ε) is referred to as the n -jet rate. This n -jet rate is explicitly defined as

$$R_n(y_{\text{cut}}) = \frac{\sigma_n(y_{\text{cut}})}{\sigma_{\text{tot}}} \equiv \frac{N_n(y_{\text{cut}})}{N_{\text{tot}}}, \quad (1)$$

where σ_n is the cross-section for the production of a hadronic event with n jets at fixed y_{cut} , σ_{tot} is the total hadronic cross-section, $N_n(y_{\text{cut}})$ is the number of events in a sample with n jets for a given value of y_{cut} and N_{tot} is the total number of events in that sample.

The differential n -jet rate was also determined. It is the derivative of the n -jet rate with respect to y_{cut} ,

$$D_n(y_{\text{cut}}) = \frac{dR_n(y_{\text{cut}})}{dy_{\text{cut}}}. \quad (2)$$

For the case when $n = 2$, the differential 2-jet rate reduces to $D_2 = y_{23}$, where y_{23} is the value of the jet resolution parameter where the event flips from a 2- to a 3-jet event. When the Durham algorithm is used to define jets the value of D_2 (denoted y_{23}^D) is also an event shape variable.

The average number of jets per event in a given sample, as a function of the jet resolution parameter, is defined to be

$$\begin{aligned} \langle N \rangle(y_{\text{cut}}) &= \frac{1}{\sigma_{\text{tot}}} \sum_n n \sigma_n(y_{\text{cut}}) \\ &= \frac{1}{N_{\text{tot}}} \sum_n n N_n(y_{\text{cut}}). \end{aligned} \quad (3)$$

A QCD prediction which matches an $\mathcal{O}(\alpha_s^2)$ (next-leading order) prediction, based on the QCD matrix elements [26], with a resummed, next-leading logarithmic

approximation (NLLA) [27] prediction, such that terms that appear in both predictions are not double counted, was fitted to data. In this analysis we used the $\ln R$ matched D_2 [22,28,29] and $\langle N \rangle$ [6,30] predictions to fit to the distributions of the observables³. The differential and average jet rates were determined using the Durham and Cambridge clustering algorithms, since resummed predictions only exist for these algorithms. This provided four separate observables (D_2^D , D_2^C , $\langle N \rangle^D$ and $\langle N \rangle^C$) which were used to determine a value of α_s at the four different centre-of-mass energy values.

5 Analysis procedure

5.1 Selection method

5.1.1 Preselection

All events within a dataset were required to contain information from both the central jet chamber and the electromagnetic calorimeter, meaning both these subdetectors must have been flagged as being on and in good operational condition. In addition events were required to be tagged as multihadronic in order to be preselected for analysis. Multihadronic events were identified using the criteria described in [31] for events with $\sqrt{s} > M_Z$ and in [32] for Z-calibration events. To pass the preselection, an event was required to contain at least seven good tracks to reduce potential backgrounds arising from the production of τ leptons ($e^+e^- \rightarrow \tau^+\tau^-$) decaying into hadrons and from two-photon interactions producing quarks. Good tracks were defined as those which had

- at least 40 hits in the jet chamber
- at least 150 MeV/c transverse momentum relative to the beam axis.
- the distance of closest approach to the interaction point in the $r - \phi$ plane satisfying $d_0 \leq 2$ cm
- the point of closest approach ≤ 25 cm from the interaction point in the z -direction

Clusters of energy in the calorimeters were also used in the analysis; good clusters were defined as those which produced a signal in at least one block in the barrel electromagnetic calorimeter corresponding to an uncorrected energy of at least 100 MeV or of 2 blocks in the endcap electromagnetic calorimeter corresponding to an uncorrected energy of 250 MeV. The hadron calorimeter was not used in this analysis.

All of the good quality tracks and clusters in the event were used to define “objects” representing particles using an algorithm (MT) to correct for double counting of energy.

³ It should be noted that the resummed theoretical calculations apply for massless quarks, while the quarks in the Monte Carlo models do have masses. No attempt was made to correct for this. In previous OPAL papers, a systematic error was estimated for this effect, and proved to be smaller than the other hadronization errors. For consistency with other LEP experiments, we now neglect this.

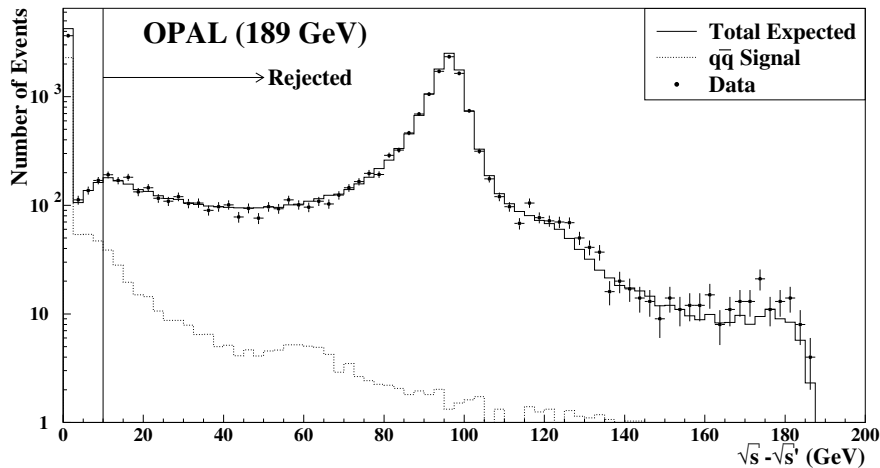


Fig. 1. The $\sqrt{s} - \sqrt{s'}$ distribution for the 189 GeV dataset. The vertical line indicates where the cut was applied. The expected non-radiative $q\bar{q}$ signal ($\sqrt{s} - \sqrt{s'_{\text{true}}} < 1$ GeV) was determined from Monte Carlo and normalized to the luminosity of the measured sample. Vertical error bars indicate the size of statistical errors, while horizontal error bars correspond to the bin width

This MT algorithm produced a uniquely defined array of track and cluster objects. The trajectories of the tracks measured in the central tracking chambers were extrapolated to the clusters in the electromagnetic calorimeters. If the energy of the cluster was less than expected from the track, then the cluster was omitted to avoid double counting of energy, since the momentum resolution for tracks was typically better than the calorimeter energy resolution. If the energy of the cluster was larger than expected the energy of the cluster was reduced by the expected amount with the remaining energy interpreted as due to photons or neutral hadrons. These remaining clusters and those which were not matched defined the four-vectors of “neutral” particles. In all cases tracks were treated as charged pions and neutral particles were treated as being massless.

5.1.2 Containment

We ensured that most particles in the event were well contained in the detector and not lost down the beam line by imposing a cut on the direction of the thrust axis [33],

$$- |\cos \theta_T| < 0.9,$$

where θ_T is the angle between the beam axis and the direction of the thrust axis. The thrust axis direction was determined from all tracks and clusters in the event, without correcting for double counting with the MT algorithm.

5.1.3 Initial state radiation (ISR) cuts

The events of interest for this analysis were $e^+e^- \rightarrow q\bar{q}$ events where the final-state $q\bar{q}$ pair had the full centre-of-mass energy. The effective centre-of-mass energy of the e^+e^- collision can be reduced by the emission of one or more ISR photons. At LEP2, approximately three quarters of the multihadronic events were “radiative return events”, where the invariant mass of the $q\bar{q}$ pair was close to the Z mass. The effective centre-of-mass energy of the collision after ISR, $\sqrt{s'}$ [34], was evaluated, and the requirement

$$- \sqrt{s} - \sqrt{s'} < 10 \text{ GeV}$$

was imposed to select full-energy events.

To calculate $\sqrt{s'}$, all isolated photon candidates with energies greater than 10 GeV were identified. The Durham jet reconstruction algorithm [31] was then used to group the remaining tracks and clusters into jets. ISR photons are often emitted close to the beam direction. Three kinematic fits were performed, under the assumptions that

- there were two undetected photons (in opposite directions along the beam pipe)
- there was one undetected photon
- all photons were observed in the detector,

respectively. The fit with the lowest number of photons which gave an acceptable χ^2 was taken [35], and $\sqrt{s'}$ was calculated from the invariant mass of the jets, excluding any photons.

The power of this cut can be seen in Fig. 1. The efficiency for selecting non-radiative $q\bar{q}$ events is given in Table 2. The purity of non-radiative events was found to be approximately 73% in all of the LEP1.5 and LEP2 data samples. Non-radiative $q\bar{q}$ events are defined as those in which $\sqrt{s} - \sqrt{s'_{\text{true}}} < 1$ GeV, where s'_{true} was determined from generator-level information in the PYTHIA samples. This ISR cut was applied to all analyzed datasets with the exception of the Z calibration data.

5.1.4 Final cuts

The dominant background to the process $e^+e^- \rightarrow Z/\gamma^* \rightarrow q\bar{q} \rightarrow \text{hadrons}$ at LEP2 came from the four-fermion process $e^+e^- \rightarrow W^+W^-$ in which one or both of the bosons decayed hadronically, producing two or four quarks in the final state. This background was expected to make up approximately 30% of all observed events which pass the first stage of cuts in each of the LEP2 datasets. These backgrounds were addressed by placing a cut [36] on two likelihood values which indicate how likely an event is to be a non-QCD four-quark or a semi-leptonic event:

- $\mathcal{L}_{q\bar{q}q\bar{q}} < 0.25$
- $\mathcal{L}_{q\bar{q}l\nu} < 0.50$

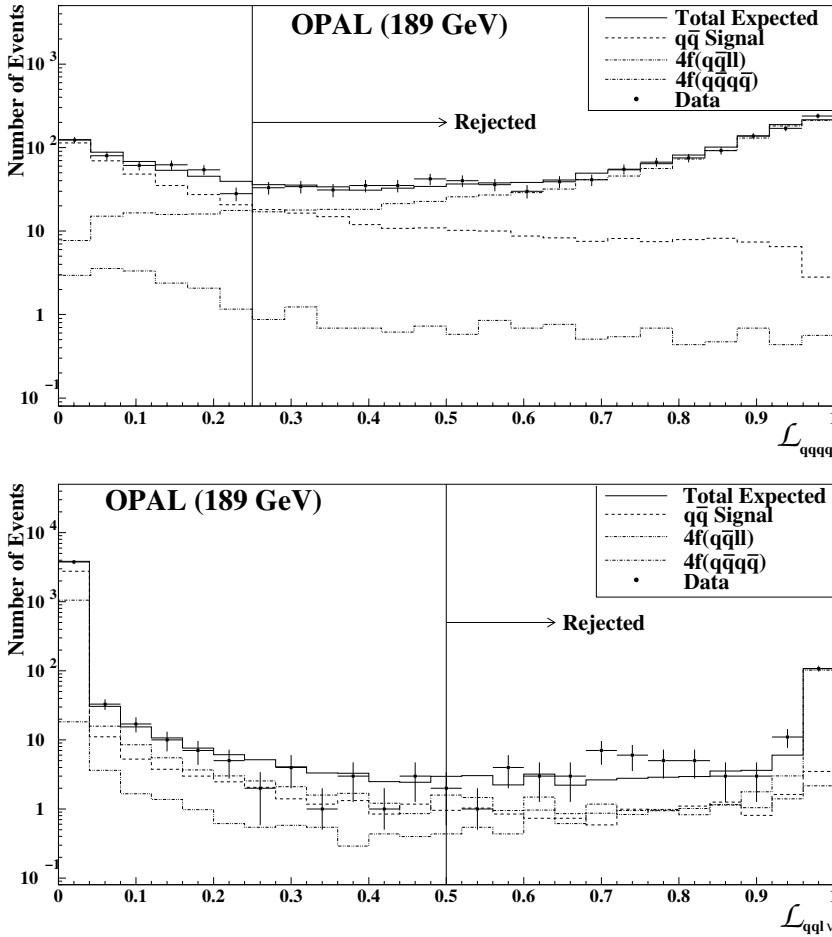


Fig. 2. Distributions of the four-quark likelihood $\mathcal{L}_{q\bar{q}q\bar{q}}$ (top) and semi-leptonic likelihood $\mathcal{L}_{q\bar{q}\ell\nu}$ (bottom) for the 189 GeV dataset. The expected contribution of the four-quark non-QCD background to the total sample is shown by the dotted line, and the expected semi-leptonic four-fermion background is represented by the dashed line. The vertical lines indicate the values of the parameters where the cuts were applied. The expected $q\bar{q}$ signal was determined from Monte Carlo and normalized to the luminosity of the measured sample. Vertical error bars indicate the size of statistical errors, while the horizontal bars correspond to the bin width

The effect of these cuts in each of the LEP2 datasets and the expected backgrounds can be seen on Fig. 2.

The four-quark likelihood value [37], $\mathcal{L}_{q\bar{q}q\bar{q}}$, was estimated from four kinematic variables describing characteristics of hadronic W^+W^- decays like their four-jet nature and angular structure. These variables were used to construct event probabilities based on two hypotheses: first, that the event was due to a hadronically decaying W^+W^- pair ($W^+W^- \rightarrow q\bar{q}q\bar{q}$) and, second, that the event was due to a hadronically decaying Z/γ^* ($e^+e^- \rightarrow Z/\gamma^* \rightarrow q\bar{q}$). The probabilities were combined to produce the discriminating likelihood, $\mathcal{L}_{q\bar{q}q\bar{q}}$. This cut reduced the expected background by approximately 80% so that it constituted only 9% of the observed number of events.

The semi-leptonic likelihood [37], $\mathcal{L}_{q\bar{q}\ell\nu}$, was based upon three separate likelihoods, one for each lepton species ($\ell = e, \mu, \tau$). Each of these likelihoods was based on ten variables describing the properties of the lepton, the jets produced by the $q\bar{q}$ pair and the missing energy carried away by the neutrino. This cut in conjunction with the cut on the four-quark likelihood removed almost 90% of the background expected in the observed LEP2 dataset. The effect of these cuts can be seen in Fig. 2.

These likelihood cuts also reduced the backgrounds arising from $e^+e^- \rightarrow ZZ$ in which one or both of the Z bosons decay hadronically. ZZ production contributed only a small

fraction of the background due to its lower cross-section compared to W^+W^- production in the energy ranges used in this experiment. The likelihood cuts were applied only to those datasets with $\sqrt{s} \geq 161$ GeV.

The expected size of the total background contribution to each dataset was determined by Monte Carlo predictions after scaling to the luminosity of the dataset. The effect of the final cuts and the expected four-fermion backgrounds for each centre-of-mass energy dataset can be seen in Table 2. As seen in the table the likelihood cuts greatly increased the purity of selected non-radiative $q\bar{q}$ events. The LEP2 datasets data with $\sqrt{s} \geq 183$ GeV were typically $\sim 70\%$ pure following the ISR cuts; however, after the final cuts this increased to a purity of 94–95%.

5.2 Monte Carlo corrections

The values of the variables R_n , D_n and $\langle N \rangle$ were determined for each accepted event using the MT-corrected tracks and clusters. These values were then compiled into histograms as a function of y_{cut} with bins of varying size. The background rejection cut did not completely remove all of the expected background events from W^+W^- and ZZ production (referred to as four-fermion background in this paper). The remaining backgrounds, taken from the Monte Carlo, were subtracted from the corresponding measured distri-

Table 2. Effect of selection cuts (see text) on the number of observed events and the expectation from Monte Carlo simulations for the $q\bar{q}$ signal plus the four-fermion background. The quoted errors are purely statistical and do not reflect any experimental or hadronization systematics. The efficiency to select non-radiative $q\bar{q}$ events with $\sqrt{s} - \sqrt{s'} < 1$ GeV and the expected fraction of four-fermion background are also shown. The four-fermion background is negligible below 161 GeV

| | Dataset | ISR Fit | Final | | Dataset | ISR Fit | Final |
|---------|-----------------|------------|------------|---------|-----------------|----------|----------|
| 91 GeV | Data | 395695 | 395695 | 192 GeV | Data | 717 | 514 |
| 130 GeV | Data | 318 | 318 | | MC Expected | 696±14 | 471±4 |
| | MC Expected | 368 ± 3 | 368 ± 3 | | non-rad eff(%) | 82.8±0.7 | 77.4±0.7 |
| | non-rad eff(%) | 85.4 ± 1.4 | 85.4 ± 1.4 | | 4f bkg frac (%) | 31.3±1.5 | 5.3±1.2 |
| 136 GeV | Data | 312 | 312 | 196 GeV | Data | 1732 | 1137 |
| | MC Expected | 329 ± 3 | 329 ± 3 | | MC Expected | 1746±24 | 1162±9 |
| | non-rad eff(%) | 85.5 ± 1.4 | 85.5 ± 1.4 | | non-rad eff(%) | 82.7±0.7 | 77.1±0.6 |
| 161 GeV | Data | 304 | 281 | | 4f bkg frac (%) | 32.6±1.0 | 5.8±1.0 |
| | MC Expected | 299 ± 4 | 274 ± 3 | 200 GeV | Data | 1636 | 1090 |
| | non-rad eff(%) | 83.4 ± 1.0 | 80.0 ± 1.0 | | MC Expected | 1717 ±25 | 1130±10 |
| | 4f bkg frac (%) | 6.2 ± 1.7 | 2.1 ± 1.4 | | non-rad eff(%) | 82.6±0.7 | 76.9±0.7 |
| 172 GeV | Data | 280 | 218 | | 4f bkg frac (%) | 33.6±1.0 | 6.0±1.0 |
| | MC Expected | 288 ± 7 | 232 ± 3 | 202 GeV | Data | 806 | 519 |
| | non-rad eff(%) | 83.3 ± 1.0 | 79.6 ± 1.0 | | MC Expected | 804±16 | 527±5 |
| | 4f bkg frac (%) | 19.7 ± 2.2 | 4.4 ± 1.7 | | non-rad eff(%) | 85.4±0.7 | 76.9±0.7 |
| 183 GeV | Data | 1456 | 1077 | | 4f bkg frac (%) | 34.1±1.4 | 6.2±1.2 |
| | MC Expected | 1478 ± 22 | 1083 ± 11 | 205 GeV | Data | 1687 | 1130 |
| | non-rad eff(%) | 83.0 ± 1.0 | 79.0 ± 1.0 | | MC Expected | 1693±25 | 1089±9 |
| | 4f bkg frac (%) | 27.3 ± 1.2 | 5.5 ± 1.2 | | non-rad eff(%) | 82.4±0.7 | 76.3±0.7 |
| 189 GeV | Data | 4448 | 3086 | | 4f bkg frac (%) | 34.8±1.0 | 6.2±1.0 |
| | MC Expected | 4497 ± 37 | 3130 ±16 | 207 GeV | Data | 2713 | 1717 |
| | non-rad eff(%) | 83.0 ± 0.4 | 78.1 ± 0.4 | | MC Expected | 2807±32 | 1804±12 |
| | 4f bkg frac (%) | 30.0 ± 0.6 | 5.5 ± 0.6 | | non-rad eff(%) | 82.7±0.5 | 76.5±0.5 |
| | | | | | 4f bkg frac (%) | 34.7±0.8 | 6.2±0.6 |

butions on a bin-by-bin basis. Systematic uncertainties in this procedure will be discussed in Sect. 6.

Corrections to the distributions were also made for effects arising from finite detector resolution and a limited detector acceptance (recall that the cut on $|\cos\theta_T|$ reduced the fiducial volume) and for residual ISR events which were not removed by the $\sqrt{s'}$ cut. These corrections were done separately for each variable and were accomplished by comparing distributions from two separate Monte Carlo samples, one of which had gone through a full detector simulation, including effects of detector resolution and acceptance and initial state radiation, called the “detector” level. The other sample used only the generator-level hadrons and had not gone through the detailed detector simulation. In this sample all short-lived particles ($\tau \leq 3 \times 10^{-10}$ s) had decayed and a requirement that $\sqrt{s} - \sqrt{s'_{\text{true}}} < 1$ GeV was imposed. This “hadron level” sample was thus expected to produce distributions arising solely from the properties of the underlying hadrons, free of any detector biases de-

termined over the full acceptance without any limitations arising from limited resolution.

Correction factors for each bin of the distributions were determined from the ratio of the two Monte Carlo distributions. Thus, any bin, i , of the measured distributions was corrected via

$$\mathcal{H}_i^{\text{data}} = \left(\frac{\mathcal{H}_i^{\text{MC}}}{\mathcal{D}_i^{\text{MC}}} \right) (\mathcal{D}_i^{\text{data}} - \mathcal{D}_i^{\text{bkgd}}) \quad (4)$$

where \mathcal{H} and \mathcal{D} represent distributions at the hadron level and the detector level respectively, and $\mathcal{D}_i^{\text{bkgd}}$ corresponds to the expected size of the total background in bin i . The hadron level was used in this analysis when determining jet rate distributions; theoretical predictions which were fitted to these distributions were obtained from computations valid at the “parton level”. The parton level corresponds to distributions that would be produced if only the partons created immediately following the e^+e^- annihilation and before the hadronization phase were used in the analysis. The parton level Monte Carlo sample was built from

quarks and gluons that were produced during the parton shower simulated by the generator before the hadronization phase began. As in the case of the hadron level, the parton level sample gave rise to distributions that were free of initial state radiation without any detector simulation applied. The correction factor determined from the ratio of the hadron level to the parton level was applied to the theoretical predictions before the fitting procedure. As will be mentioned later (Sect. 7.2.1), it was verified that these hadronic corrections did not exceed 10% over the fitted range of the variable. Plots, illustrating the size of this correction can be found in our earlier publication [6]. This factor corrected the prediction to the hadron level so that it could be compared to the corrected hadron level distribution determined from the data,

$$\mathcal{H}_i^{\text{pred}} = \left(\frac{\mathcal{H}_i^{\text{MC}}}{\mathcal{P}_i^{\text{MC}}} \right) \mathcal{P}_i^{\text{pred}} \quad (5)$$

where \mathcal{H} and \mathcal{P} represent distributions at the hadron level and the parton level respectively. Statistical errors on the data distributions included the effects of Monte Carlo statistics.

6 Systematic uncertainties

6.1 Experimental systematic variations

Several selection algorithms and selection cuts were varied to determine their impact on the results of the analysis. In all cases the result from the variation was compared to the result from the standard selection, and the difference was taken as a contribution to the total systematic error. The systematic variations that were used (given in descending order of their contributions to the overall experimental uncertainty) are:

Track-Cluster Matching Systematic errors relevant to the definition of objects defined within the tracking chambers and calorimeters, and hence used for jet clustering, were estimated by comparing the results using the MT package to a method using all selected tracks and clusters without taking into account the possibility of double counting.

Detector Correction The uncertainty in modelling the detector was investigated by using HERWIG Monte Carlo datasets in place of PYTHIA to correct from the detector level to the hadron level. The hadronization correction was still performed using the PYTHIA Monte Carlo samples.

Containment The constraint on the direction of the thrust axis was tightened to $|\cos \theta_T| < 0.7$, restricting events to the barrel region.

ISR A possible systematic effect introduced through the selection of events with little or no initial state radiation was estimated by repeating the selection using a second method of ISR determination [38]. This alternative method assumed while performing the kinematic fit that there was always a single photon which either

escaped undetected down the beam line or was detected in the electromagnetic calorimeter.

$\mathcal{L}_{q\bar{q}q\bar{q}}$ and $\mathcal{L}_{q\bar{q}\ell\nu}$ To account for the systematic uncertainty that arises from the value of the cut placed on the W^+W^- hadronic and semi-leptonic likelihoods, the values of the likelihood cuts were changed to 0.1 and 0.4 for $\mathcal{L}_{q\bar{q}q\bar{q}}$ and to 0.25 and 0.75 for $\mathcal{L}_{q\bar{q}\ell\nu}$. In each case the largest deviation was taken as the systematic uncertainty.

Backgrounds To account for uncertainties introduced during background subtraction, arising from imprecise knowledge of the four-fermion cross-sections, these cross-sections were conservatively varied by $\pm 5\%$ and the largest deviation from the standard value was used to determine an overall systematic error.

It should be noted that there was no single dominant contribution to the overall experimental systematic error. In general, the largest contributions occurred for the track-cluster matching and detector correction variations, while the least significant contributions came from the variations of the background cross-section and the semi-leptonic likelihood.

6.2 Hadronization systematic variations

Systematic uncertainties arise from the modelling of the hadronization. These were estimated by using HERWIG Monte Carlo samples, which employ a different hadronization model, in the determination of the theoretical prediction correction factor (5). The PYTHIA-HERWIG differences were taken to be the systematic uncertainties from hadronization. The statistical component due to limited Monte Carlo statistics was included in the total determination of the hadronization uncertainty.

6.3 Theoretical systematic variations

Three further systematic variations were considered when fitting the predictions to the data to determine α_s .

Fit Range We investigated the choice of the range of bins used in performing the fit of the theoretical prediction to the data. The fit range was first increased by two bins, by adding one bin to each endpoint of the fit range. In the case where one of the endpoints was already at the maximum allowed bin value, only the other point was extended by one bin. A second variation decreased the fit range by two bins, by removing a bin from each endpoint of the fit range. The largest deviation from the standard fit value was taken as the contribution to the systematic error.

Renormalization Scale The second fit-related systematic variation accounted for the uncertainty due to the dependence of the fixed order and resummed predictions on the renormalization scale, x_μ , where $x_\mu = \mu/\sqrt{s}$. The value of x_μ , which was set to 1.0 in the standard fits, was varied to 0.5 and 2.0 respectively. The largest deviation from the standard fit value was taken as the

Table 3. Summary of systematic variations applied to the datasets. The * indicates which of the systematic variations were used for the 91 GeV dataset

| Category | Error Source | Standard | Variation 1 | Variation 2 |
|----------|--|------------|-------------------|-------------|
| | Track-Cluster Matching* | MT objects | Tracks + Clusters | |
| | Detector Correction* | PYTHIA | HERWIG | |
| | Containment ($ \cos\theta_T $)* | < 0.9 | < 0.7 | |
| (exp.) | qqqq ($\mathcal{L}_{q\bar{q}q\bar{q}}$) | < 0.25 | < 0.1 | < 0.4 |
| | qq ν ($\mathcal{L}_{q\bar{q}\ell\nu}$) | < 0.5 | < 0.75 | < 0.25 |
| | ISR Algorithm | Ref. [31] | Ref. [38] | |
| | Backgrounds (σ_{bkgd}) | 1.0 | +5% | -5% |
| (had.) | Hadronization* | PYTHIA | HERWIG | |
| (theo.) | Fit Range* | min 6 bins | +2(1) bins | -2 bins |
| | Renorm. Scale Dep. (x_μ)* | 1 | 0.5 | 2.0 |

contribution to the systematic error representing renormalization scale uncertainty.

Logarithm rescaling, x_L In the resummation process for event shape variables, like D_2 , there is an arbitrariness due to the definition of the logarithms which were resummed. In this analysis we used $\alpha_s \ln(1/y_{\text{cut}})$; however, this could be generalized to powers of $\alpha_s \ln[1/(x_L y_{\text{cut}})]$ [39]. The standard value for this rescaling ($x_L=1$) was varied to $\frac{4}{9}$ and $\frac{9}{4}$, to investigate the systematic effect of this arbitrariness. No analogous rescaling prescription has yet been developed for the case of the average jet rates, so the x_L variations are only shown for the D_2 distributions for comparison to the x_μ variation. Hence this variation was not used in the determination of the total theoretical systematic error and was included only as a cross-check and for comparison with other analyses. A comprehensive study of the combination of various theoretical variations to produce a global theoretical uncertainty for event shape distributions is given in [40].

The differences between the standard result and those due to the above variations were separated into three categories: experimental, hadronization and theory (see Table 3). The variations in each of these categories were added in quadrature and the result taken as the systematic uncertainty for that category. In the case of asymmetric errors, the error was symmetrized by taking the largest systematic variation and applying it as the full systematic contribution. A summary of the systematic variations is provided in Table 3.

7 Results

Data from thirteen datasets were used in this analysis: one Z-calibration dataset, two LEP1.5 datasets and ten datasets from LEP2. These thirteen datasets were combined to produce four higher statistics datasets ($\sqrt{s} = M_Z$, 133, 177 and 197 GeV) which were analysed separately. The raw distributions (n -jet fractions, D_2 and $\langle N \rangle$) for each of the datasets were determined as functions of the

jet resolution parameters defined using four different jet clustering algorithms. The distributions underwent a bin-by-bin correction which included subtraction of expected backgrounds and correction for detector and residual ISR effects. Systematic effects were examined by varying the parameters used in selecting events (see Table 3). The difference between the corrected distributions using these variations and those from the standard selection then determined the size of the systematic error on each bin of the distribution.

Matched predictions of NLO and resummed calculations were fitted to the corrected Durham and Cambridge D_2 and $\langle N \rangle$ distributions over a predetermined fit range (see Sect. 7.2.1), taking into account bin-to-bin correlations in $\langle N \rangle$ (see Sect. 7.2.2). These fits provided four values of α_s with statistical and systematic errors at each of the four centre-of-mass energies. Taking into account the statistical and systematic correlations between the four measurements, they were combined into a single α_s result at each energy.

7.1 n -Jet Fractions

The n -jet fractions for the Z-calibration sample and those for the LEP1.5 and the two LEP2 samples are shown in Figs. 3 to 7. Each plot shows the fraction of events in a given sample that were determined to be n -jets for a given value of the jet resolution parameter at the hadron level⁴. The jet fractions were calculated using four different clustering algorithms. For the Cone algorithm, results from using both the R and ε resolution parameters are plotted, showing the individual n -jet fractions for $n \leq 2$, $n = 3$ and $n \geq 4$ in Figs. 3 to 4. Similarly the n -jet fractions for the JADE algorithm are shown in Fig. 5 for $n = 2, 3, 4, 5$. Figures 6 to 7 show the individual n -jet fractions for $n = 2, 3, 4, 5$ for the Durham and Cambridge algorithms respectively. The error bars on the points represent the total statistical and systematic errors added in quadrature. The Monte Carlo

⁴ Further details of the data will be made available in the HEPDATA database, <http://durpdg.dur.ac.uk/HEPDATA>.

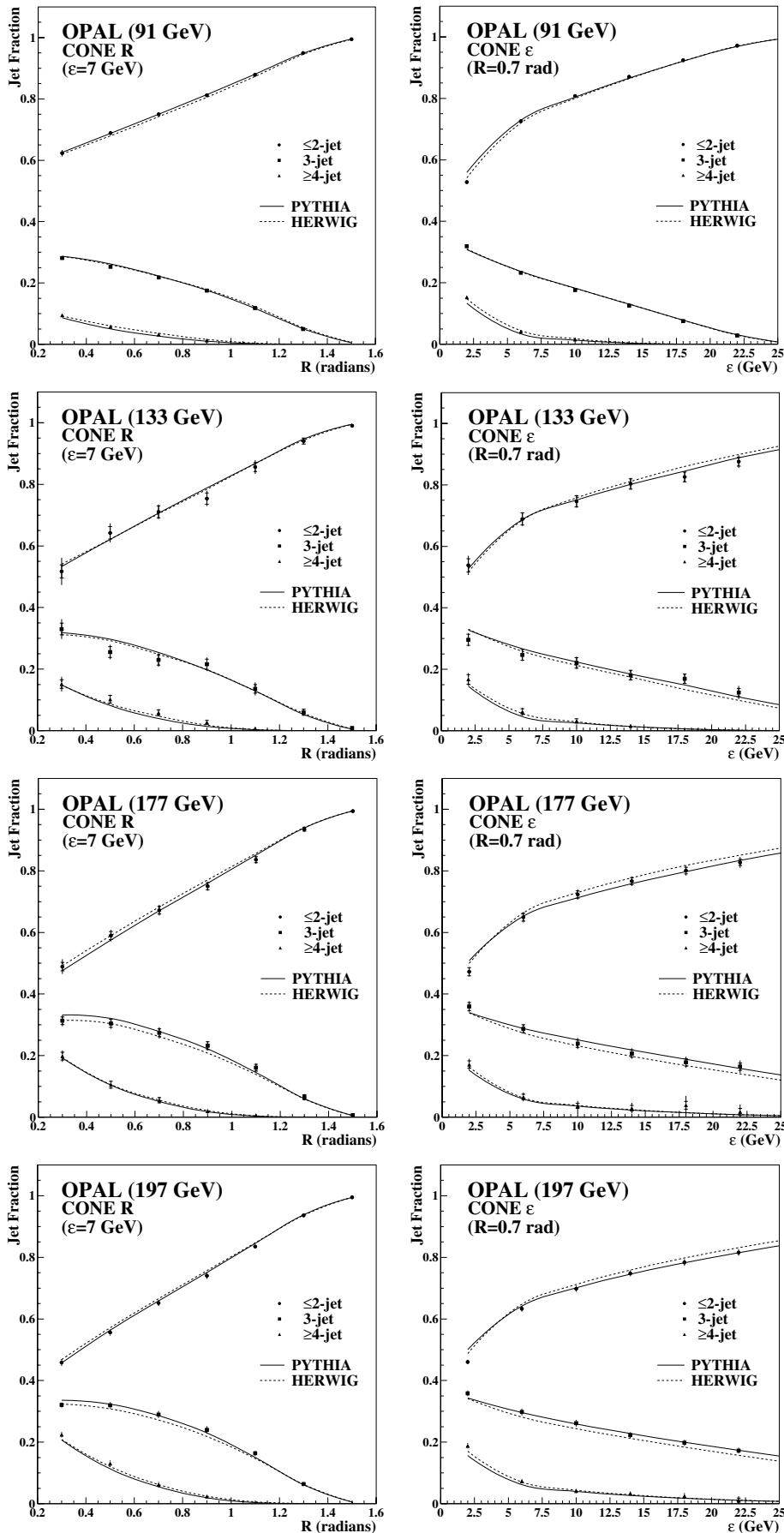


Fig. 3. The hadron level n -jet rates for the R and ϵ variants of the Cone algorithm for the data with $\sqrt{s} = 91$ GeV (top) and $\sqrt{s} = 133$ GeV (bottom). In all plots the JETSET/PYTHIA and HERWIG Monte Carlo expectations are represented by the curves. Outer error bars indicate total errors while the inner bars indicate statistical errors

Fig. 4. The hadron level n -jet rates for the R and ϵ variants of the Cone algorithm for the data with $\sqrt{s} = 179$ GeV (top) and $\sqrt{s} = 198$ GeV (bottom). In all plots the PYTHIA and HERWIG Monte Carlo expectations are represented by the curves. Outer error bars indicate total errors while the inner bars indicate statistical errors

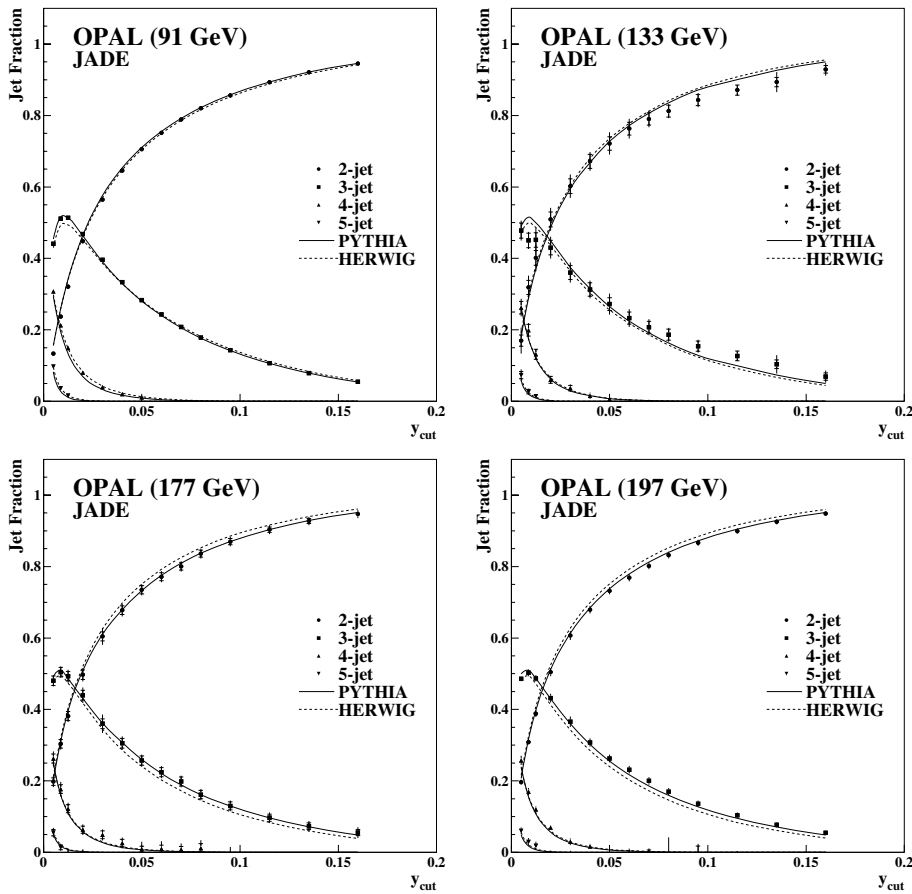


Fig. 5. The hadron level n -jet rates for the JADE algorithm for the data with $\sqrt{s} = 91$ GeV (top left), $\sqrt{s} = 133$ GeV (top right), $\sqrt{s} = 177$ GeV (bottom left) and $\sqrt{s} = 197$ GeV (bottom right). In all plots the JETSET/PYTHIA and HERWIG Monte Carlo expectations are represented by the curves. Outer error bars indicate total errors while the inner bars indicate statistical errors

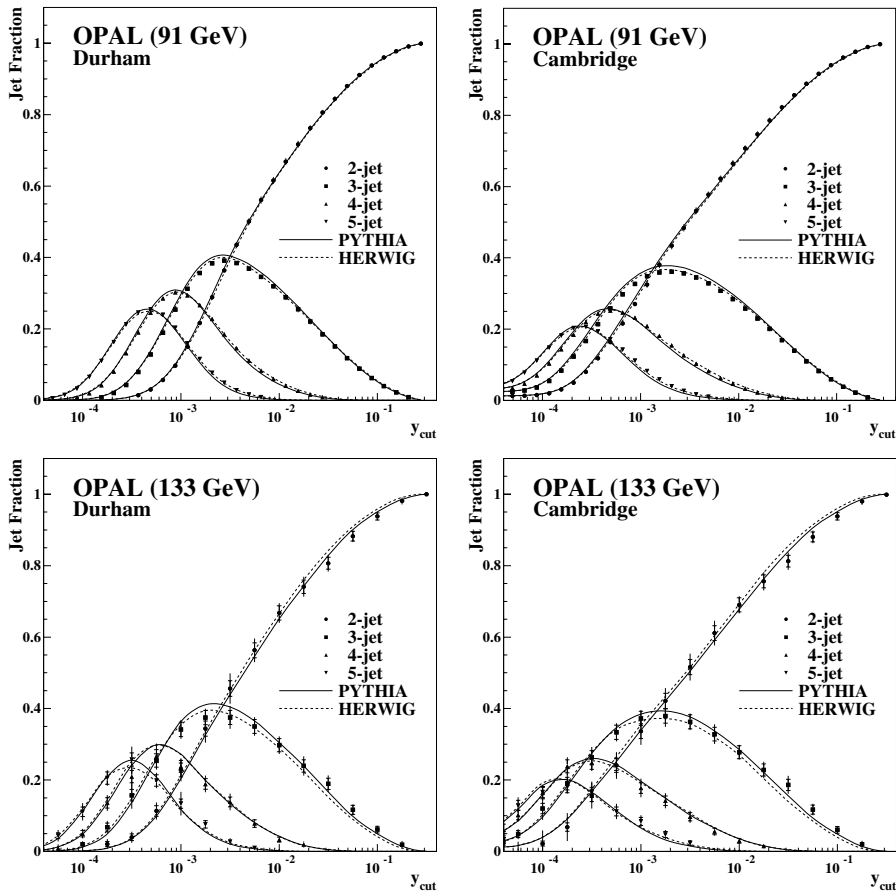


Fig. 6. The hadron level n -jet rates for the Durham and Cambridge algorithms for the data with $\sqrt{s} = 91$ GeV (top) and $\sqrt{s} = 133$ GeV (bottom). In all plots the JETSET/PYTHIA and HERWIG Monte Carlo expectations are represented by the curves. Outer error bars indicate total errors while the inner bars indicate statistical errors

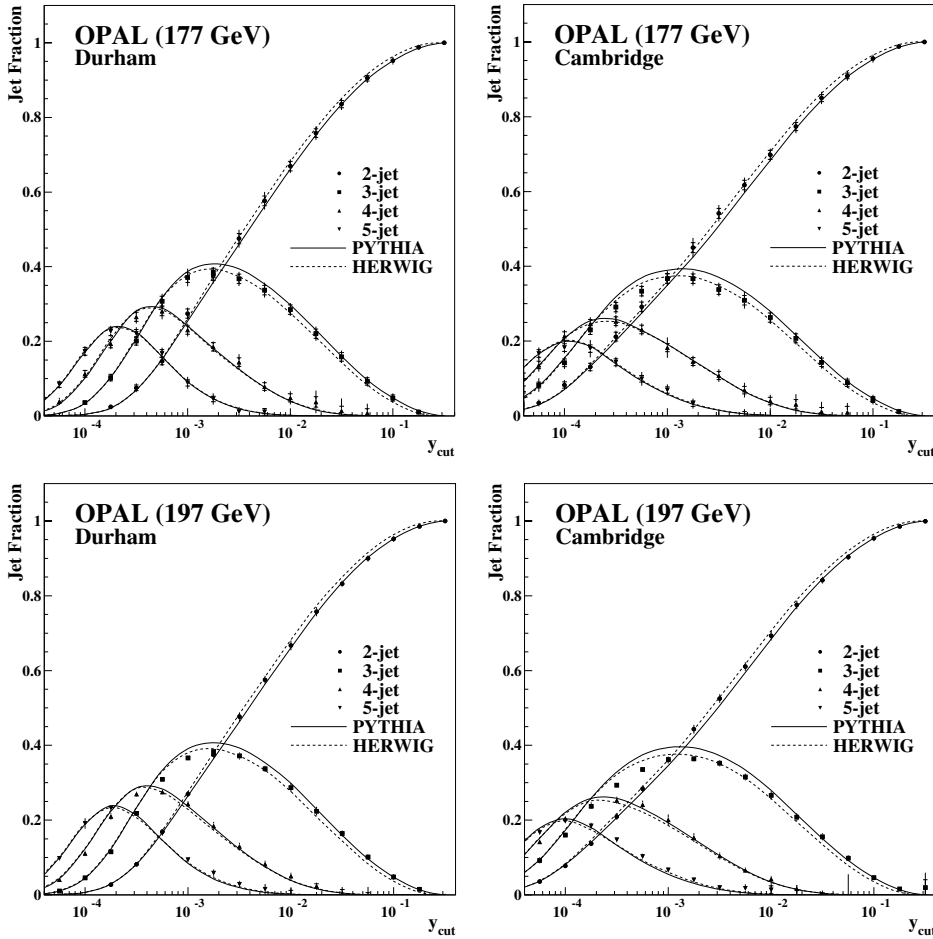


Fig. 7. The hadron level n -jet rates for the Durham and Cambridge algorithms for the data with $\sqrt{s} = 179$ GeV (top) and $\sqrt{s} = 198$ GeV (bottom). In all plots the PYTHIA and HERWIG Monte Carlo expectations are represented by the curves. Outer error bars indicate total errors while the inner bars indicate statistical errors

expectations corresponding to PYTHIA and HERWIG are also displayed on each plot, for each algorithm and energy. The Monte Carlo expectations match the measured n -jet fractions reasonably well.

The differential two-, three- and four-jet rates, D_n , are plotted as a function of y_{cut} for the Durham and Cambridge algorithms for the Z-calibration, the LEP1.5 and the two LEP2 data samples in Figs. 8 and 9 respectively. The bin sizes of the Z-calibration data are roughly comparable to the expected experimental resolution and therefore some positive bin-to-bin correlations may be present. All sources of correlation are taken into account in our fitting procedure discussed below. The average jet rates are plotted as a function of y_{cut} for the Durham and Cambridge algorithms for the Z-calibration, the LEP1.5 and the two LEP2 data samples in Figs. 10 and 11 respectively. The curves on all the plots represent the expected Monte Carlo distributions. There is good agreement between the data and the expectations from both PYTHIA and HERWIG.

7.2 Fits to determine the value of α_s

7.2.1 Differential two-jet rates

The range over which the D_2 and $\langle N \rangle$ distributions were fitted was determined by splitting the 91 GeV (and 189 GeV)

PYTHIA Monte Carlo samples into 100 statistically independent subsamples. The QCD predictions for each distribution were fitted to each of these subsamples, with α_s as a free parameter, for a number of possible end-points of the fit range (requiring that the fit range be at least six bins). A χ^2 per degree of freedom was determined for each fit range. The χ^2 values were averaged over the 100 subsamples. The fit range that produced the smallest average χ^2 per degree of freedom was chosen to be the default fit range for the Z-calibration (and high energy) datasets. Where no clear χ^2 minimum was found, the largest reasonable range was chosen. The size of the fit range was then adjusted to ensure the range did not extend into a region where the hadronization corrections exceeded 10%, in particular at smaller y_{cut} values. A further adjustment was made to exclude fit ranges where one of the endpoints produced a contribution of more than 30% to the overall χ^2 value. Potential correlations introduced in the correction of the distributions were accounted by including a covariance matrix in the χ^2 fit. The covariance matrices were determined in the manner detailed in Sect. 7.2.2.

The fits of the $\ln R$ matching prediction to the D_2 distribution are shown in Fig. 12 for the Cambridge algorithm and in Fig. 13 for the Durham algorithm. The values of α_s determined from the $\ln R$ fits for the four datasets, together with a complete breakdown of statistical and systematic

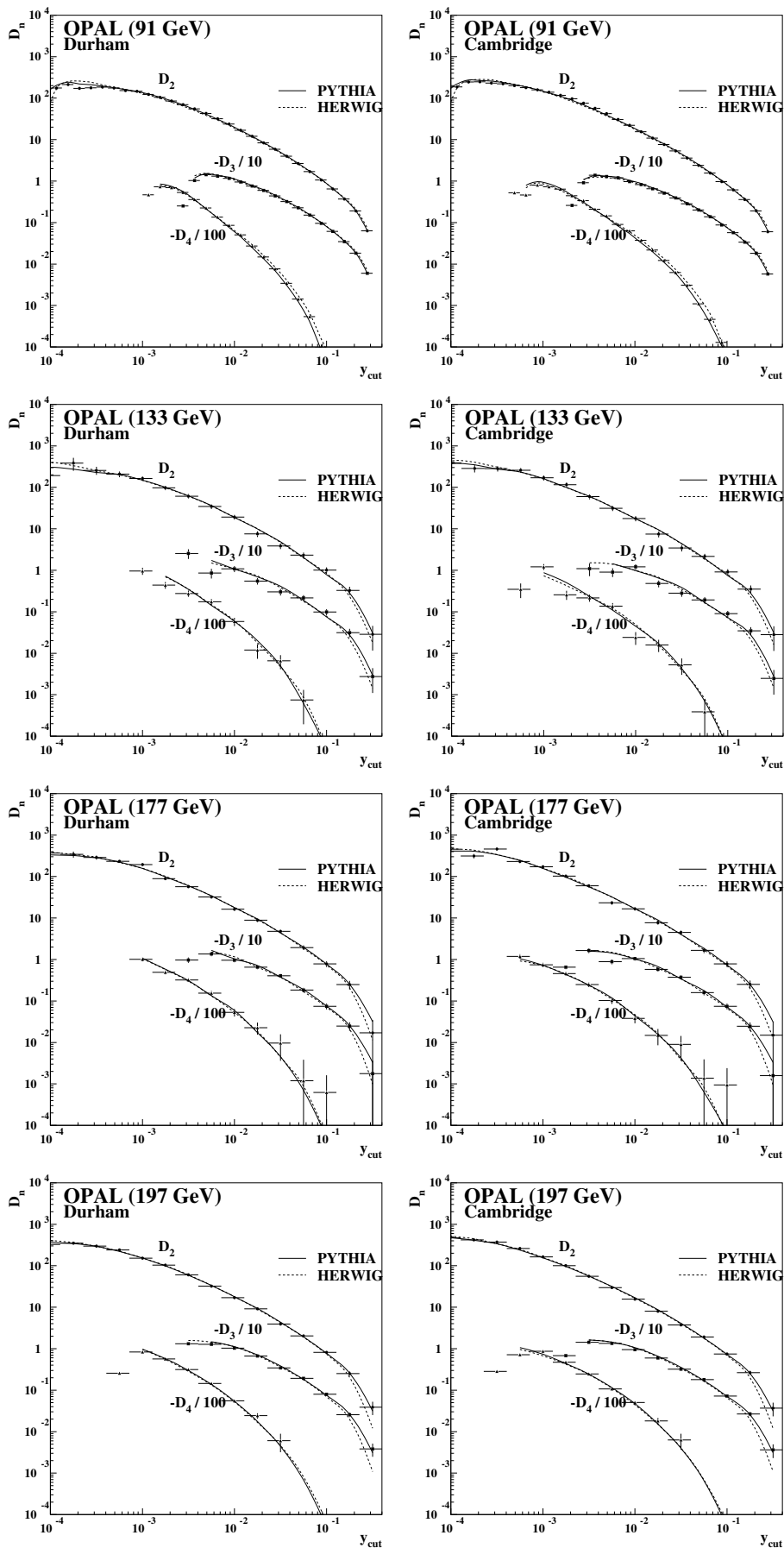


Fig. 8. The differential 2-, 3- and 4-jet rates at the hadron level as a function of y_{cut} for the Cambridge and Durham algorithms for data with $\sqrt{s} = 91$ GeV (top) and for the LEP1.5 combined dataset at $\sqrt{s} = 133$ GeV (bottom). PYTHIA and HERWIG Monte Carlo expectations are represented by the curves. The differential 3- and 4-jet rates have a negative slope in the region of large y_{cut} and therefore the negative values of D_3 and D_4 are plotted (the positive values are not seen on the curves). Note that D_3 and D_4 are scaled down by one and two orders of magnitude, respectively, for clarity. Error bars indicate total (statistical + systematic) errors

Fig. 9. The differential 2-, 3- and 4-jet rates at the hadron level as a function of y_{cut} for the Cambridge and Durham algorithms for two combined LEP2 datasets with $\sqrt{s} = 177$ GeV (top) and $\sqrt{s} = 197$ GeV (bottom). PYTHIA and HERWIG Monte Carlo expectations are represented by the curves. The differential 3- and 4-jet rates have a negative slope in the region of large y_{cut} and therefore the negative values of D_3 and D_4 are plotted (the positive values are not seen on the curves). Note that D_3 and D_4 are scaled down by one and two orders of magnitude, respectively, for clarity. Error bars indicate total (statistical + systematic) errors

Table 4. Determination of α_s and the breakdown of statistical and systematic errors from the fit to the Cambridge differential two-jet rate distribution (D_2^C) for all centre-of-mass energy values. The quality of the fit is characterized by the chi-square (χ^2) and the number of degrees of freedom (dof). The theoretical error includes contributions from the fit range and x_μ variations and excludes the x_L variation, which is used only for comparison

| | 91 GeV | 133 GeV | 177 GeV | 197 GeV |
|---|--------------|--------------|--------------|--------------|
| α_s (\sqrt{s}) | 0.1147 | 0.1071 | 0.0991 | 0.0996 |
| Fit Range ($-\log_{10} y_{\text{cut}}$) | 2.06–0.81 | 2.75–0.75 | 2.75–0.75 | 2.75–0.75 |
| χ^2 /dof | 7.98/10 | 7.74/8 | 10.33/8 | 8.63/8 |
| Experimental | ± 0.0017 | ± 0.0032 | ± 0.0019 | ± 0.0009 |
| Hadronization | ± 0.0026 | ± 0.0007 | ± 0.0005 | ± 0.0005 |
| Fit Range Variation | ± 0.0005 | ± 0.0035 | ± 0.0008 | ± 0.0004 |
| x_μ $x_\mu=2.0$ | +0.0030 | +0.0030 | +0.0023 | +0.0020 |
| x_μ $x_\mu=0.5$ | -0.0010 | -0.0014 | -0.0011 | -0.0010 |
| x_L $x_L=9/4$ | (+0.0040) | (+0.0033) | (+0.0026) | (+0.0025) |
| x_L $x_L=4/9$ | (-0.0026) | (-0.0022) | (-0.0018) | (-0.0019) |
| Theoretical | ± 0.0030 | ± 0.0046 | ± 0.0024 | ± 0.0020 |
| Total Stat. | ± 0.0004 | ± 0.0026 | ± 0.0019 | ± 0.0008 |
| Total Syst. | ± 0.0043 | ± 0.0056 | ± 0.0031 | ± 0.0022 |
| Total Error | ± 0.0043 | ± 0.0062 | ± 0.0037 | ± 0.0024 |

Table 5. Determination of α_s and the breakdown of statistical and systematic errors from the fit to the Durham differential two-jet rate distribution (D_2^D) for all centre-of-mass energy values. The quality of the fit is characterized by the chi-square (χ^2) and the number of degrees of freedom (dof). The theoretical error includes contributions from the fit range and x_μ variations and excludes the x_L variation, which is used only for comparison

| | 91 GeV | 133 GeV | 177 GeV | 197 GeV |
|---|--------------|--------------|--------------|--------------|
| α_s (\sqrt{s}) | 0.1199 | 0.1129 | 0.1060 | 0.1064 |
| Fit Range ($-\log_{10} y_{\text{cut}}$) | 1.81–0.68 | 2.75–0.75 | 2.75–0.75 | 2.50–0.75 |
| χ^2 /dof | 4.02/9 | 7.02/8 | 1.59/8 | 9.10/7 |
| Experimental | ± 0.0025 | ± 0.0026 | ± 0.0013 | ± 0.0007 |
| Hadronization | ± 0.0017 | ± 0.0005 | ± 0.0005 | ± 0.0007 |
| Fit Range Variation | ± 0.0005 | ± 0.0020 | ± 0.0035 | ± 0.0015 |
| x_μ $x_\mu=2.0$ | +0.0037 | +0.0028 | +0.0031 | +0.0030 |
| x_μ $x_\mu=0.5$ | -0.0013 | -0.0020 | -0.0017 | -0.0014 |
| x_L $x_L=9/4$ | (+0.0044) | (+0.0028) | (+0.0031) | (+0.0032) |
| x_L $x_L=4/9$ | (-0.0027) | (-0.0026) | (-0.0022) | (-0.0020) |
| Theoretical | ± 0.0037 | ± 0.0034 | ± 0.0047 | ± 0.0034 |
| Total Stat. | ± 0.0004 | ± 0.0026 | ± 0.0019 | ± 0.0009 |
| Total Syst. | ± 0.0048 | ± 0.0043 | ± 0.0049 | ± 0.0035 |
| Total Error | ± 0.0048 | ± 0.0050 | ± 0.0052 | ± 0.0037 |

uncertainties, are given in Tables 4 and 5 for the Cambridge and Durham⁵ algorithms respectively.

⁵ Note that [7] also determines α_s using D_2^D , denoted there as y_{23}^D . The small differences between the results have been investigated in detail, and are not significant. They may be attributed to differences in fit regions, the use of statistically different Monte Carlo samples, and the adoption of slightly different strategies for the assessment of theoretical errors.

7.2.2 Average jet rates

For the average jet rates, $\langle N \rangle$, the statistical errors were strongly correlated between points, since the same events were used to determine $\langle N \rangle$ at each value of y_{cut} . The fit was performed using a correlated χ^2 fit in which the covariance matrix was determined from the PYTHIA Monte Carlo sample divided into many detector level subsamples. Each subsample was then corrected to the hadron level using a second, statistically independent, PYTHIA sample.

Table 6. Determination of α_s and the breakdown of statistical and systematic errors from the fit to the Cambridge average jet rate distribution ($\langle N \rangle^C$) for all centre-of-mass energy values. The quality of the fit is characterized by the chi-square (χ^2) and the number of degrees of freedom (dof)

| | 91 GeV | 133 GeV | 177 GeV | 197 GeV |
|---|----------------------------|------------------------|------------------------|------------------------|
| $\alpha_s (\sqrt{s})$ | 0.1254 | 0.1158 | 0.1064 | 0.1066 |
| Fit Range ($-\log_{10} y_{\text{cut}}$) | 2.56–0.56 | 2.75–0.50 | 2.75–0.50 | 2.75–0.50 |
| χ^2/dof | 20.70/16 | 7.38/9 | 2.57/9 | 5.16/9 |
| Experimental | ± 0.0013 | ± 0.0040 | ± 0.0030 | ± 0.0015 |
| Hadronization | ± 0.0033 | ± 0.0034 | ± 0.0018 | ± 0.0017 |
| Fit Range Variation | ± 0.0008 | ± 0.0042 | ± 0.0020 | ± 0.0010 |
| x_μ | $x_\mu=2.0$ $x_\mu=0.5$ | $+0.0036$ -0.0001 | $+0.0021$ -0.0001 | $+0.0016$ -0.0001 |
| Theoretical | ± 0.0037 | ± 0.0047 | ± 0.0026 | ± 0.0019 |
| Total Stat. | ± 0.0005 | ± 0.0033 | ± 0.0023 | ± 0.0009 |
| Total Syst. | ± 0.0051 | ± 0.0070 | ± 0.0042 | ± 0.0030 |
| Total Error | ± 0.0051 | ± 0.0077 | ± 0.0048 | ± 0.0031 |

Table 7. Determination of α_s and the breakdown of statistical and systematic errors from the fit to the Durham average jet rate distribution ($\langle N \rangle^D$) for all centre-of-mass energy values. The quality of the fit is characterized by the chi-square (χ^2) and the number of degrees of freedom (dof)

| | 91 GeV | 133 GeV | 177 GeV | 197 GeV |
|---|----------------------------|------------------------|------------------------|------------------------|
| $\alpha_s (\sqrt{s})$ | 0.1272 | 0.1193 | 0.1103 | 0.1106 |
| Fit Range ($-\log_{10} y_{\text{cut}}$) | 2.56–0.56 | 2.75–0.50 | 2.75–0.50 | 2.75–0.50 |
| χ^2/dof | 7.90/16 | 4.73/9 | 0.81/9 | 4.14/9 |
| Experimental | ± 0.0006 | ± 0.0039 | ± 0.0026 | ± 0.0013 |
| Hadronization | ± 0.0039 | ± 0.0037 | ± 0.0009 | ± 0.0007 |
| Fit Range Variation | ± 0.0003 | ± 0.0029 | ± 0.0003 | ± 0.0008 |
| x_μ | $x_\mu=2.0$ $x_\mu=0.5$ | $+0.0034$ -0.0013 | $+0.0028$ -0.0010 | $+0.0023$ -0.0008 |
| Theoretical | ± 0.0034 | ± 0.0040 | ± 0.0023 | ± 0.0024 |
| Total Stat. | ± 0.0004 | ± 0.0030 | ± 0.0021 | ± 0.0008 |
| Total Syst. | ± 0.0052 | ± 0.0067 | ± 0.0035 | ± 0.0028 |
| Total Error | ± 0.0052 | ± 0.0073 | ± 0.0041 | ± 0.0029 |

There were 1500 subsamples created for the Z-calibration dataset and 1000 subsamples created for the high energy datasets. The number of subsamples was chosen so that the elements of the covariance matrix would be stable and would have negligible fluctuations. These subsamples were used to build a standard covariance matrix, which was then converted into a correlation matrix. The statistical errors on the bins of the data distributions were then applied to this correlation matrix to produce the covariance matrix used in the fits.

An example of the fit of the $\ln R$ matching prediction to the $\langle N \rangle$ distribution is seen in Fig. 14 for the Cambridge algorithm and in Fig. 15 for the Durham algorithm. The values of α_s determined from the $\ln R$ fits for the four datasets, together with a complete breakdown of statistical and systematic uncertainties, are given in Tables 6 and 7 for the Cambridge and Durham algorithms respectively.

7.2.3 Running of α_s

The four values of α_s determined from the D_2 and $\langle N \rangle$ distributions were combined into a single value at each centre-of-mass energy. The large statistical correlations between the four values were handled in a manner similar to that for the bin-to-bin correlations of the average jet rates. For each of the four distributions, 1000 Monte Carlo samples were used to determine the statistical correlations between the α_s values. The correlation matrices determined for the four separate energy points are given in Table 8. Using the statistical error for the α_s value from each observable, the statistical covariance matrix was then determined. The full covariance matrix also included contributions from experimental, hadronization and theoretical uncertainties

$$V = V_{\text{stat}} + V_{\text{expt}} + V_{\text{hadr}} + V_{\text{theo}}. \quad (6)$$

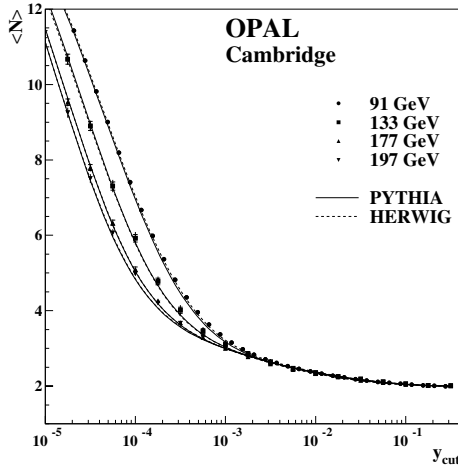


Fig. 10. The average jet rates at the hadron level as a function of y_{cut} for the Cambridge algorithm for the all centre-of-mass energies. PYTHIA and HERWIG Monte Carlo expectations are represented by the curves. Error bars indicate total (statistical + systematic) errors

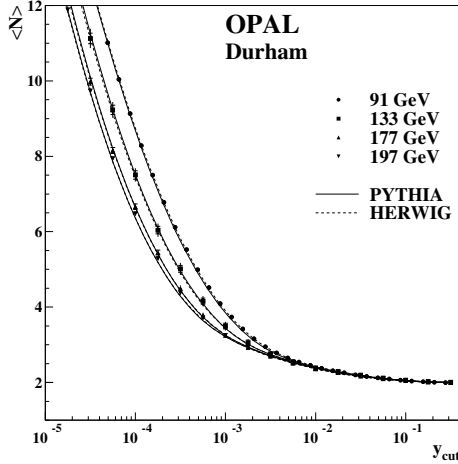


Fig. 11. The average jet rates at the hadron level as a function of y_{cut} for the Durham algorithm for the all centre-of-mass energies. PYTHIA and HERWIG Monte Carlo expectations are represented by the curves. Error bars indicate total (statistical + systematic) errors

A weight was determined for each of the α_s values from the inverse of the covariance matrix, $w_i = \sum_j (V^{-1})_{ij} / \sum_{ij} (V^{-1})_{ij}$. The combined α_s value was then determined from the weighted sum of the α_s values. Only statistical and experimental systematic uncertainties were allowed to contribute to the off diagonal elements of the covariance matrix V , to ensure undesirable features such as negative weights were avoided. Hadronization and theoretical systematics were added only to the diagonal elements of the covariance matrix.

The statistical and experimental uncertainties on the combined α_s value were determined from the product of the weights with the individual covariance matrices [41],

$$\sigma_{\text{err}}^2 = w^T V_{\text{err}} w \quad \text{where err} = \text{stat, expt.} \quad (7)$$

The hadronization and theoretical systematic uncertainties were determined by repeating the combination for each sys-

Table 8. Statistical correlation matrix determined for the combination of α_s measurements from the four different observables for each of the four different datasets. Each matrix element is presented as a percentage (%)

| | | | | |
|-----------------------|---------|---------|-----------------------|-----------------------|
| 91 GeV | D_2^C | D_2^D | $\langle N \rangle^C$ | $\langle N \rangle^D$ |
| D_2^C | 100 | 85 | 53 | 46 |
| D_2^D | 85 | 100 | 56 | 47 |
| $\langle N \rangle^C$ | 53 | 56 | 100 | 75 |
| $\langle N \rangle^D$ | 46 | 47 | 75 | 100 |
| 133 GeV | D_2^C | D_2^D | $\langle N \rangle^C$ | $\langle N \rangle^D$ |
| D_2^C | 100 | 83 | 70 | 64 |
| D_2^D | 83 | 100 | 64 | 63 |
| $\langle N \rangle^C$ | 70 | 64 | 100 | 81 |
| $\langle N \rangle^D$ | 64 | 63 | 81 | 100 |
| 177 GeV | D_2^C | D_2^D | $\langle N \rangle^C$ | $\langle N \rangle^D$ |
| D_2^C | 100 | 84 | 73 | 66 |
| D_2^D | 84 | 100 | 71 | 72 |
| $\langle N \rangle^C$ | 73 | 71 | 100 | 82 |
| $\langle N \rangle^D$ | 66 | 72 | 82 | 100 |
| 197 GeV | D_2^C | D_2^D | $\langle N \rangle^C$ | $\langle N \rangle^D$ |
| D_2^C | 100 | 90 | 77 | 68 |
| D_2^D | 90 | 100 | 72 | 73 |
| $\langle N \rangle^C$ | 77 | 72 | 100 | 82 |
| $\langle N \rangle^D$ | 68 | 73 | 82 | 100 |

tematic variation separately using the same weights. The difference between these systematic combinations and the central value is taken as the systematic contribution to the error on the central value. For the experimental systematic covariance matrix, off diagonal elements were determined using the “minimum overlap” method, that involves taking the covariance between any pair of systematic error contributions to be equal to the smaller of the two variances,

$$(V_{\text{expt}})_{ij} = \min[(V_{\text{expt}})_{ii}, (V_{\text{expt}})_{jj}]. \quad (8)$$

The values of α_s determined for each centre-of-mass energy are given in Table 9 along with the breakdown of the uncertainties, both statistical and systematic. A comparison of the combined α_s values in Table 9 with those determined from the individual D_2 and $\langle N \rangle$ distributions is seen in Fig. 16. Taking the combined α_s values at each centre-of-mass energy as an input, the value of α_s can be run back to the Z pole using an $\mathcal{O}(\alpha_s^3)$ prediction. These $\alpha_s(M_Z)$ values are also shown in Table 9 and plotted against the world average value of $\alpha_s(M_Z) = 0.1187 \pm 0.0020$ [42] in Fig. 17. Using these values and taking into account proper statistical and systematic correlations, a weighted mean of $\alpha_s(M_Z) = 0.1177 \pm 0.0006(\text{stat.}) \pm 0.0012(\text{expt.}) \pm 0.0010(\text{had.}) \pm 0.0032(\text{theo.})$ is determined. The four combined α_s values are plotted in Fig. 18 as a function of the centre-of-mass energy, com-

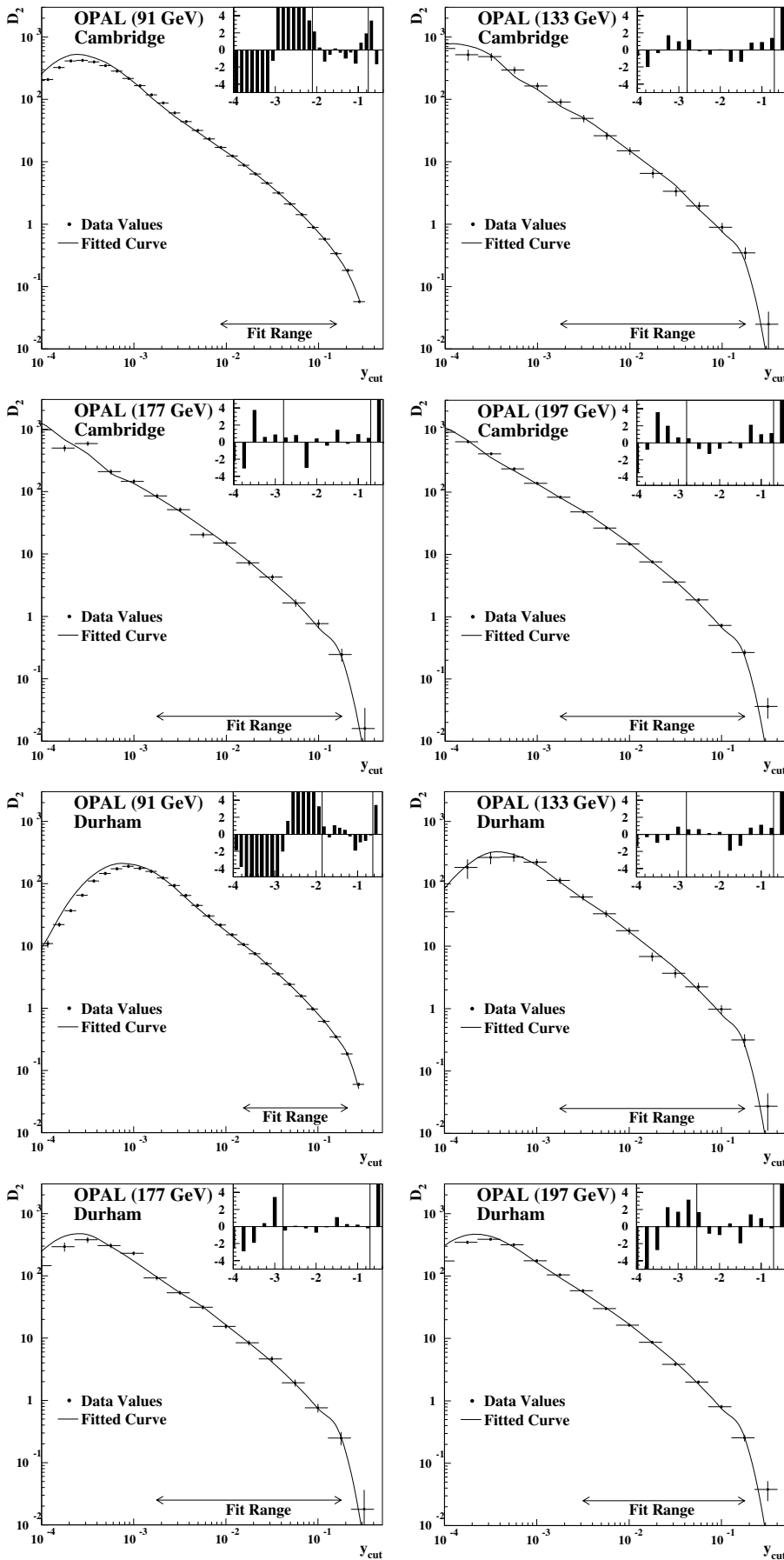


Fig. 12. Fits of the parton level D_2 distribution using the Cambridge algorithm as a function of y_{cut} to the $\ln R$ prediction for the 91 GeV (top left), 133 GeV (top right), 179 GeV (bottom left) and 198 GeV (bottom right) datasets. The inset plots show the pull of each point from the line of best fit vs $\log_{10} y_{cut}$. The two vertical lines delineate the fit range. Vertical error bars represent statistical errors only

Fig. 13. Fits of the parton level D_2 distribution using the Durham algorithm as a function of y_{cut} to the $\ln R$ prediction for the 91 GeV (top left), 133 GeV (top right), 179 GeV (bottom left) and 198 GeV (bottom right) datasets. The inset plots show the pull of each point from the line of best fit vs $\log_{10} y_{cut}$. The two vertical lines delineate the fit range. Vertical Error bars represent statistical errors only

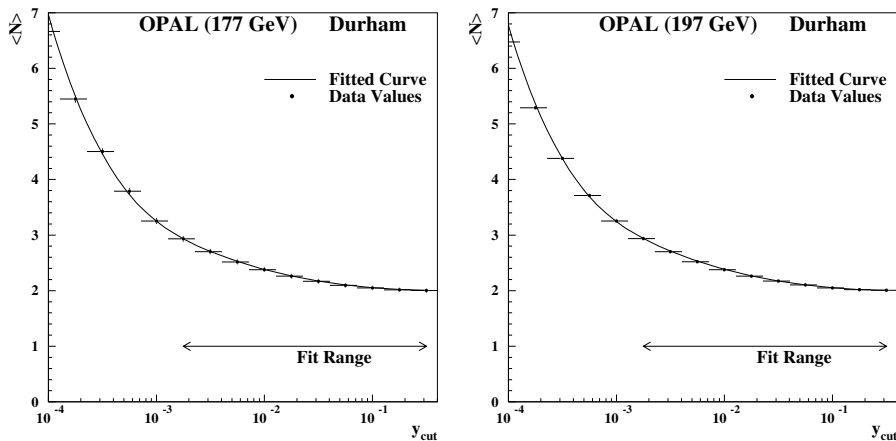
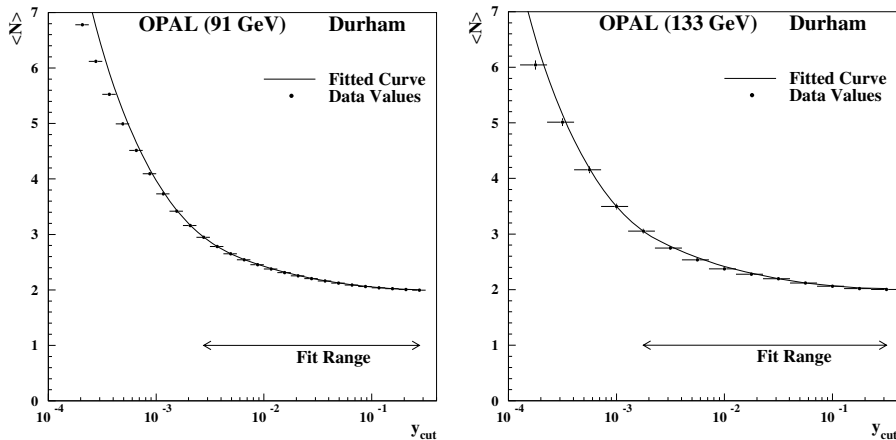
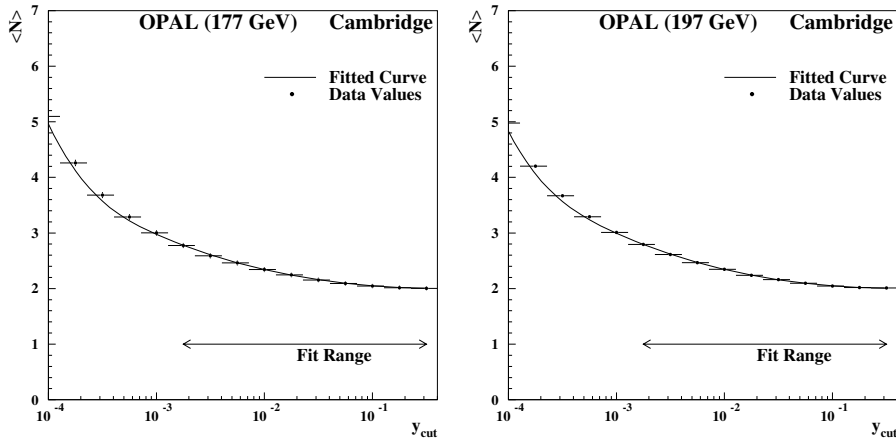
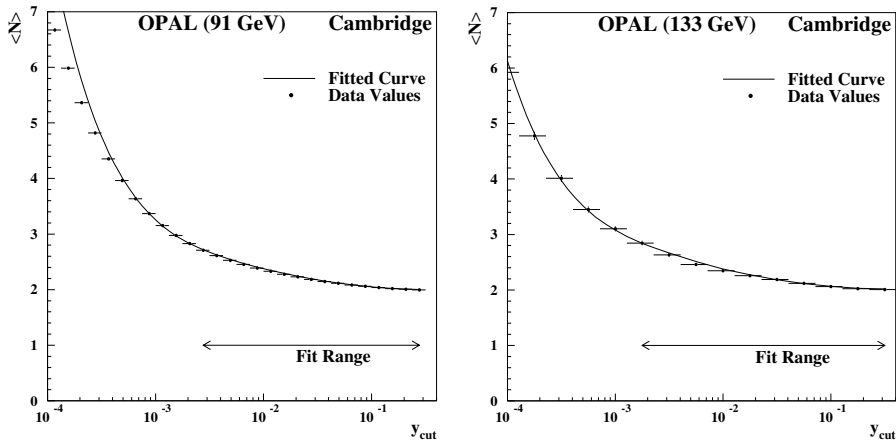


Fig. 14. Fits of the parton level $\langle N \rangle$ distribution using the Cambridge algorithm as a function of y_{cut} to the $\ln R$ prediction for the 91 GeV (top left), 133 GeV (top right), 179 GeV (bottom left) and 198 GeV (bottom right) datasets. Vertical error bars represent statistical errors only. Note: pulls are not shown as there are very large bin-to-bin correlations

Fig. 15. Fits of the parton level $\langle N \rangle$ distribution using the Durham algorithm as a function of y_{cut} to the $\ln R$ prediction for the 91 GeV (top left), 130 GeV (top right), 189 GeV (bottom left) and 207 GeV (bottom right) datasets. Vertical error bars represent statistical errors only. Note: pulls are not shown as there are very large bin-to-bin correlations

Table 9. Values of α_s determined from the weighted average of the individual α_s results at each centre-of-mass energy, along with statistical and systematic errors. The value α_s for the dataset when run back to the Z pole are given on the right half of the table with full statistical and systematic errors

| \sqrt{s} (GeV) | Value for \sqrt{s} | | | | | Value at M_Z | | | | |
|------------------|----------------------|-------------------------|------------------------|-------------------------|--------------------------|----------------|-------------------------|------------------------|-------------------------|--------------------------|
| | α_s | $\sigma_{\text{stat.}}$ | $\sigma_{\text{exp.}}$ | $\sigma_{\text{hadr.}}$ | σ_{theory} | α_s | $\sigma_{\text{stat.}}$ | $\sigma_{\text{exp.}}$ | $\sigma_{\text{hadr.}}$ | σ_{theory} |
| M_Z | 0.1213 | 0.0004 | 0.0013 | 0.0029 | 0.0034 | 0.1213 | 0.0004 | 0.0013 | 0.0029 | 0.0034 |
| 133 | 0.1126 | 0.0025 | 0.0028 | 0.0007 | 0.0039 | 0.1191 | 0.0028 | 0.0031 | 0.0008 | 0.0044 |
| 177 | 0.1039 | 0.0018 | 0.0018 | 0.0001 | 0.0028 | 0.1140 | 0.0021 | 0.0021 | 0.0001 | 0.0033 |
| 197 | 0.1046 | 0.0008 | 0.0009 | 0.0002 | 0.0023 | 0.1163 | 0.0010 | 0.0012 | 0.0003 | 0.0029 |

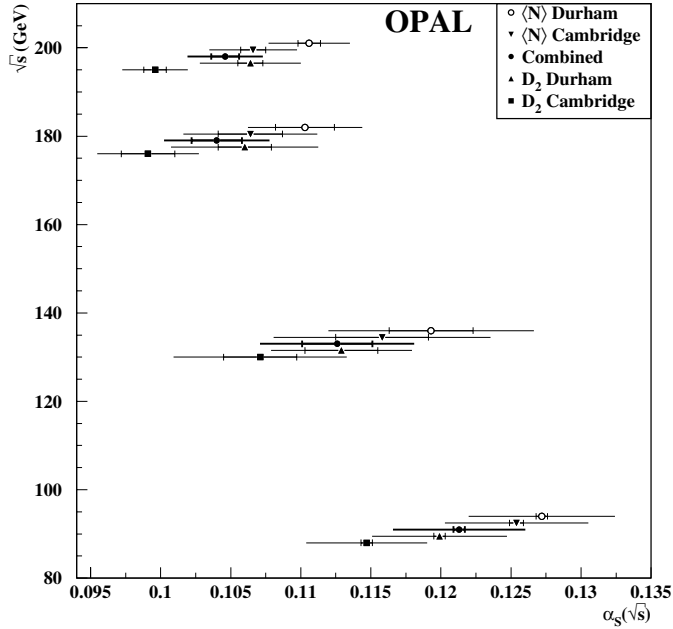


Fig. 16. Comparison of combined α_s values with those determined separately from the D_2^D , D_2^C , $\langle N \rangle^C$ and $\langle N \rangle^D$ distributions for each centre-of-mass energy dataset. Outer error bars indicate the size of the total errors while inner bars indicate the size of the statistical errors

pared to the $\mathcal{O}(\alpha_s^2)$ energy evolution of α_s based on the determined value of $\alpha_s(M_Z)$.

8 Summary and conclusion

Data from twelve LEP 1.5 and LEP 2 datasets, with centre-of-mass energies ranging from 130 GeV to 209 GeV, were combined into three higher statistics datasets. These datasets and one combining several Z-calibration runs at 91 GeV were used to determine the n -jet fractions, the differential n -jet rates and the average jet rates for each of the energies. The different jet multiplicity distributions were compared to both PYTHIA and HERWIG Monte Carlo expectations.

Hadron-level n -jet fractions were determined using four jet-clustering algorithms, Cone, JADE, Durham and Cambridge. For the Cone algorithm, measurements of the fraction of events with $n \leq 2$, 3 , ≥ 4 jets were presented as functions of R and ε . In the case of JADE, Durham and

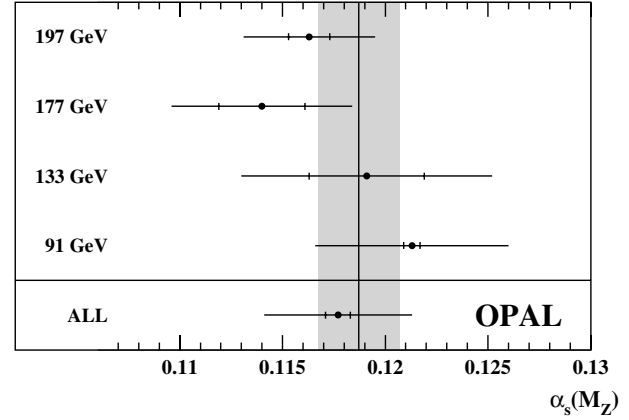


Fig. 17. Comparison of the α_s values after having been run back to the Z pole for each of the datasets. The point labelled by 'ALL' represents the value determined from the weighted mean of the four combined $\alpha_s(M_Z)$ determinations. The shaded band corresponds to the one standard deviation range of the world average value of $\alpha_s(M_Z)$ [42]. The inner error bars represent the statistical errors while the outer error bars represent the total error

Cambridge, measurements of the 2-, 3-, 4-, and 5-jet fractions were presented as functions of the jet resolution parameter, y_{cut} . In all cases there was generally good agreement between the measured jet fractions and the Monte Carlo expectations.

Hadron-level determinations of the differential n -jet and average jet rates were performed for the Durham and Cambridge algorithms. The differential two-jet rate, D_2 , and the average jet rate, $\langle N \rangle$ were used to determine the value of $\alpha_s(\sqrt{s})$ for each of the four combined energy points. The determinations were carried out by fitting the $\ln R$ matching predictions, appropriately corrected to the hadron level, to the hadron level data distributions over an appropriate range of y_{cut} , with $\alpha_s(\sqrt{s})$ as the variable parameter. The running of $\alpha_s(\sqrt{s})$ was demonstrated by comparing the four values of α_s as determined from the combined datasets as a function of their centre-of-mass energies.

Using the measured values of $\alpha_s(\sqrt{s})$ a value of $\alpha_s(M_Z)$ was determined for each of the four datasets. A weighted mean taking account of correlations determined a final value of the strong coupling at $\sqrt{s} = M_Z$ of

$$\alpha_s(M_Z) = 0.1177 \pm 0.0006(\text{stat.}) \pm 0.0012(\text{expt.}) \\ \pm 0.0010(\text{had.}) \pm 0.0032(\text{theo.})$$

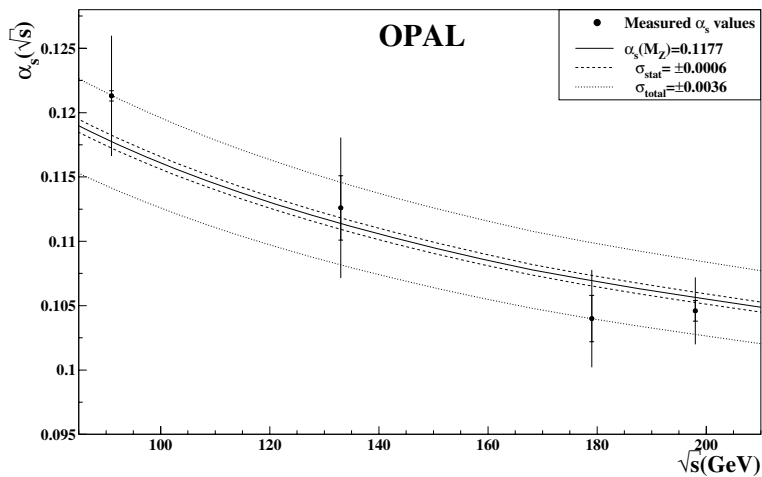


Fig. 18. The running of α_s as a function of centre-of-mass energy. The points correspond to the Z-calibration dataset, the combined LEP1.5 data and the two LEP2 regions 161–183 and 189–209, respectively. The solid line corresponds to the expectation based on the weighted mean of the four combined $\alpha_s(M_Z)$ determinations, and the outer dashed line to the total (statistical + systematic) uncertainty

where the error contains contributions from statistical, experimental, hadronization and theoretical uncertainties. The error on the determined value is slightly larger than that for the previous OPAL publication [6] which also used resummed predictions for D_2 and average jet rate distributions, but explored slightly different energy ranges, including 35 and 45 GeV, and all LEP energies up to only 189 GeV. The values of α_s reported here supersede those of our earlier publication [6]. There is good agreement between these four values of $\alpha_s(M_Z)$ measured in this analysis and previous determinations of α_s summarized in [1] and with the world average value of 0.1187 ± 0.0020 [42].

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